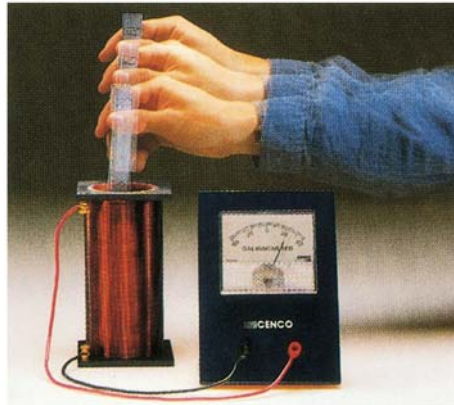


Capítulo 5

Campos electromagnéticos dependientes del tiempo



$$\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0}$$

$$\vec{\nabla} \times \vec{E} = 0$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{B} = \mu_0 \vec{J}$$

Ecuaciones de Maxwell para campos estáticos

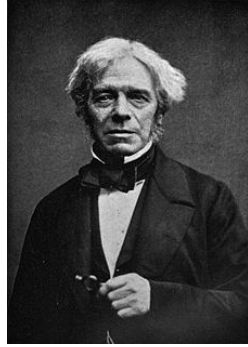
Electricidad

Magnetismo



**dos fenómenos
independientes**

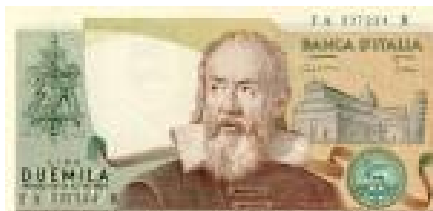
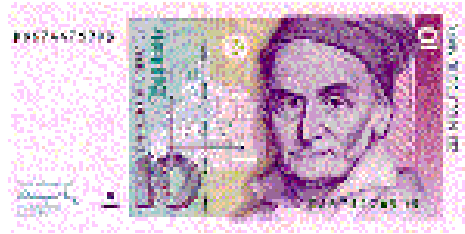
Ley de Faraday-Henry



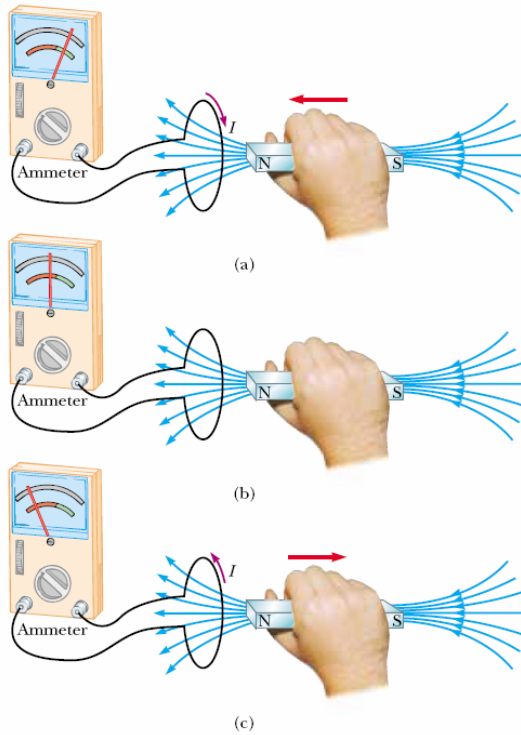
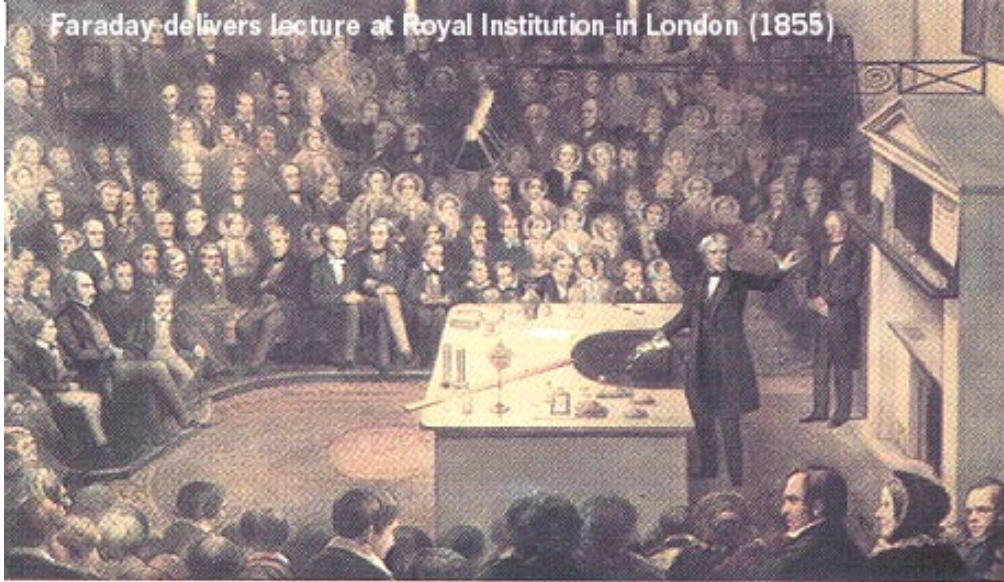
Michael Faraday (1791-1867)



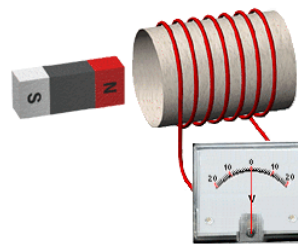
Joseph Henry (1797-1878)



Faraday delivers lecture at Royal Institution in London (1855)



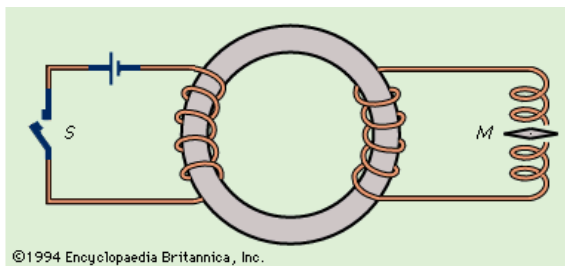
Faradays Law of Induction



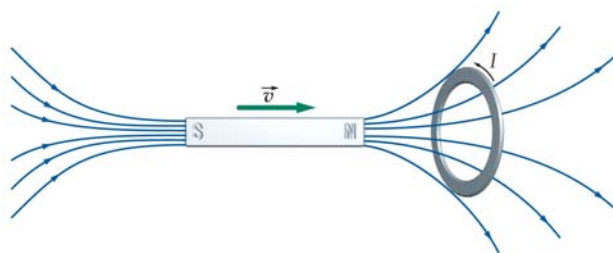
Kieran Mckenzie

Una corriente eléctrica puede producirse mediante un campo magnético variable

Experimento de Faraday

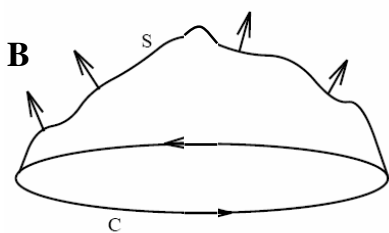


Un campo magnético variable induce una fem en el circuito secundario



$$\mathcal{E} = -\frac{d\Phi_B}{dt}$$

Enunciado general de la ley de Faraday-Henry



$$\mathcal{E} = -\frac{d\phi_m}{dt}$$

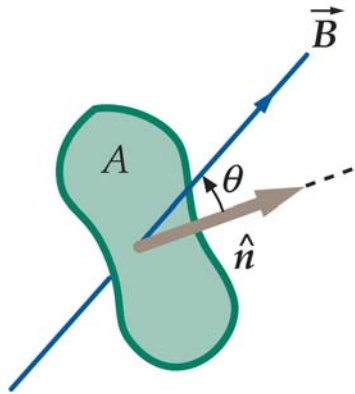
$$\phi_m = \int_S \vec{B} \cdot \hat{n} dA = \int_S B_n dA$$

$$\mathcal{E} = \oint_C \vec{E}_{nc} \cdot d\vec{\ell} = -\frac{d}{dt} \int_S \vec{B} \cdot \hat{n} dA = -\frac{d\phi_m}{dt}$$

En una bobina con N vueltas, todas de la misma área

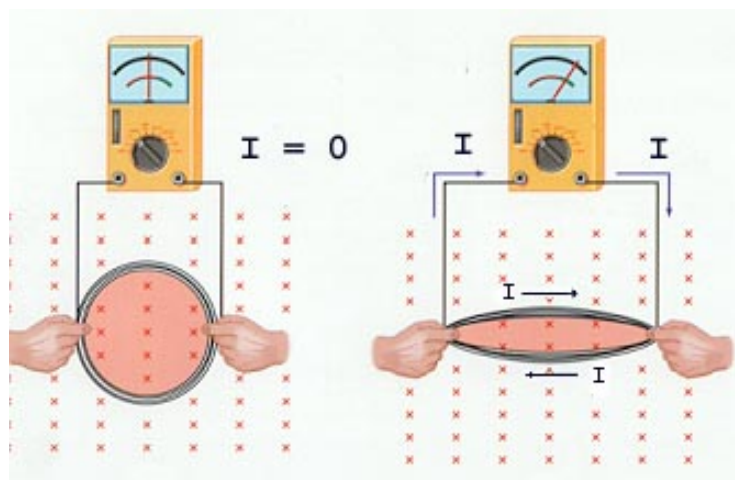
$$\mathcal{E} = -N \frac{d\Phi_B}{dt}$$

$\Phi_B(t)$?

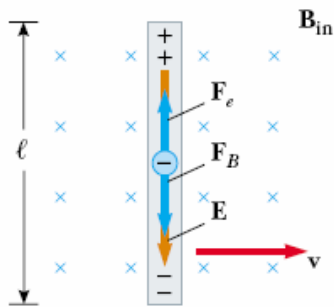
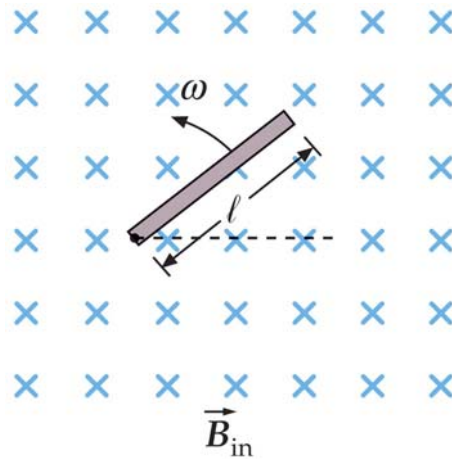


$$\mathcal{E} = -\frac{d}{dt} (BA \cos \theta)$$

- B(t)
- A(t)
- $\theta(t)$



fem de movimiento

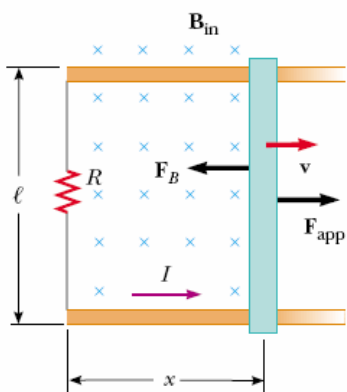


$$qE = qvB$$

$$E = vB$$

$$\Delta V = E\ell = B\ell v$$

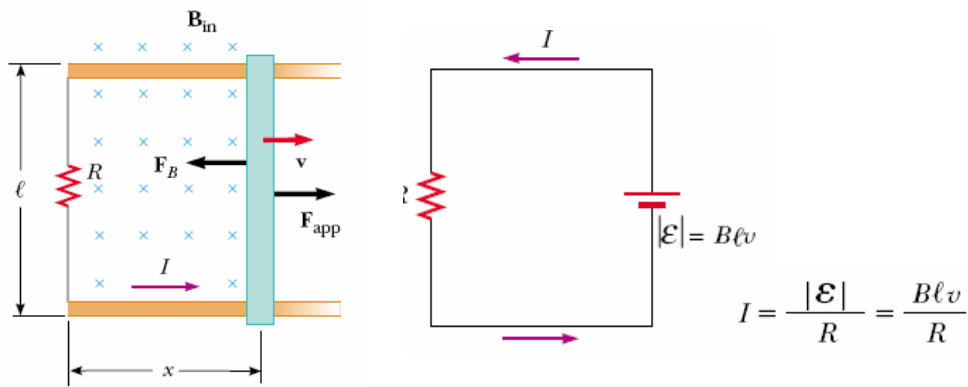
Si se invierte v , se invierte la polaridad de V



$$\Phi_B = B\ell x$$

$$\mathcal{E} = -\frac{d\Phi_B}{dt} = -\frac{d}{dt}(B\ell x) = -B\ell \frac{dx}{dt}$$

$$\mathcal{E} = -B\ell v$$



Consideración energética: si $v = \text{cte}$ $\longrightarrow F_{\text{app}} = I\ell B$

La potencia entregada por la fuerza aplicada es:

$$\mathcal{P} = F_{\text{app}}v = (I\ell B)v = \frac{B^2 \ell^2 v^2}{R} = \frac{\mathcal{E}^2}{R}$$

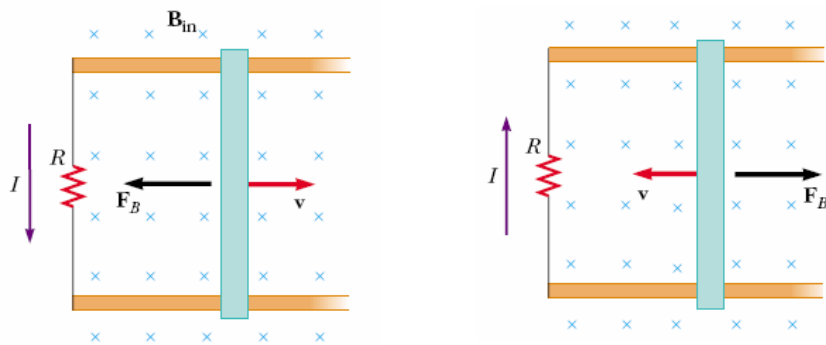
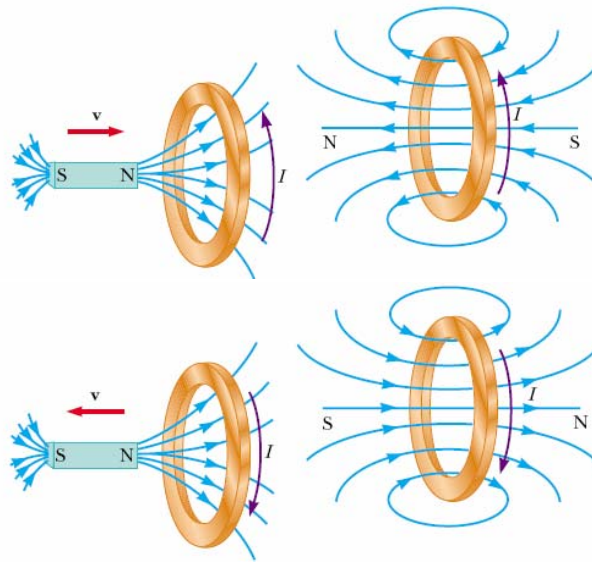
Energía Mecánica \longrightarrow **Energía Eléctrica** \longrightarrow **Energía Térmica**

Ley de Lenz



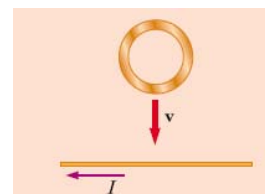
Heinrich Lenz (1804-1865)

La polaridad de una fem inducida es tal que tiende a producir una corriente eléctrica que creara un flujo magnético que se opone al cambio de Φ_B a través del lazo.

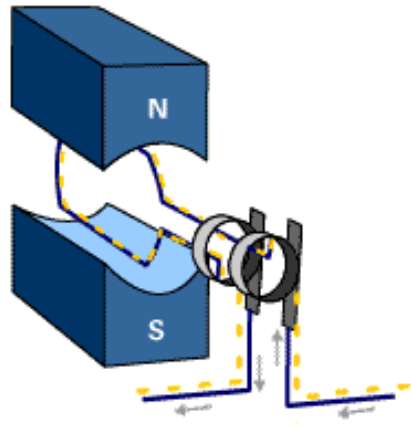


Qué sucede si el sentido de I es contrario al obtenido utilizando la ley de Lenz en los dos casos anteriores ?

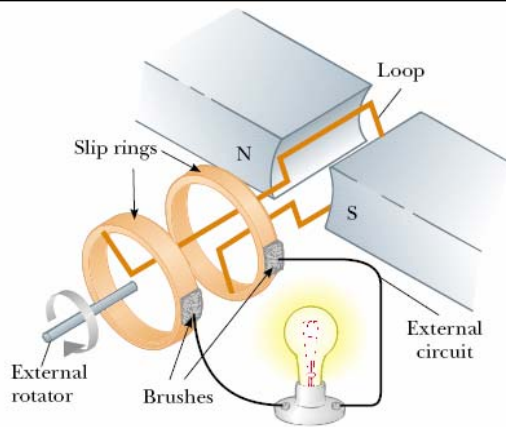
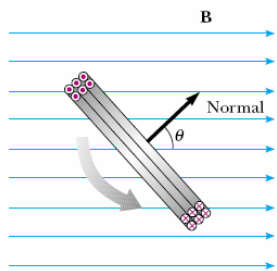
Determinar el sentido de la corriente inducida en la espira circular.



Aplicaciones de la ley de Faraday



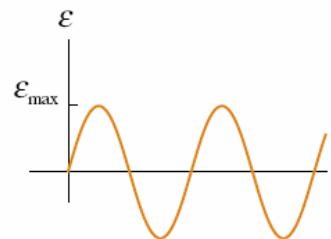
Generador eléctrico

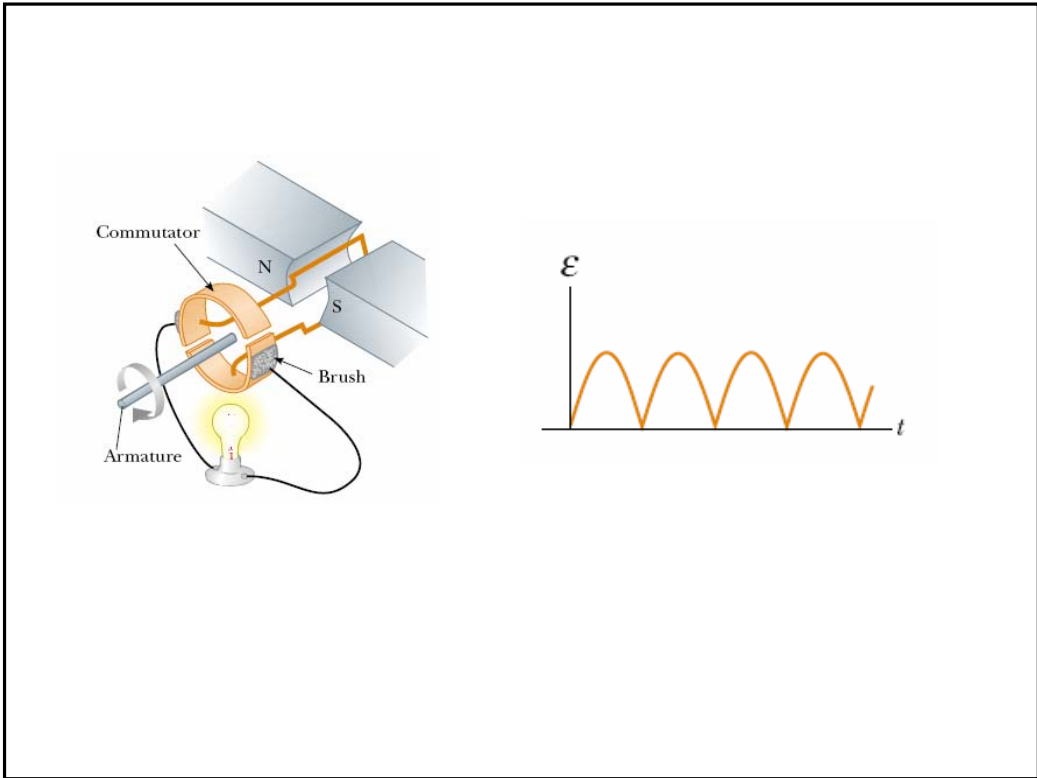


$$\Phi_B = BA \cos \theta = BA \cos \omega t$$

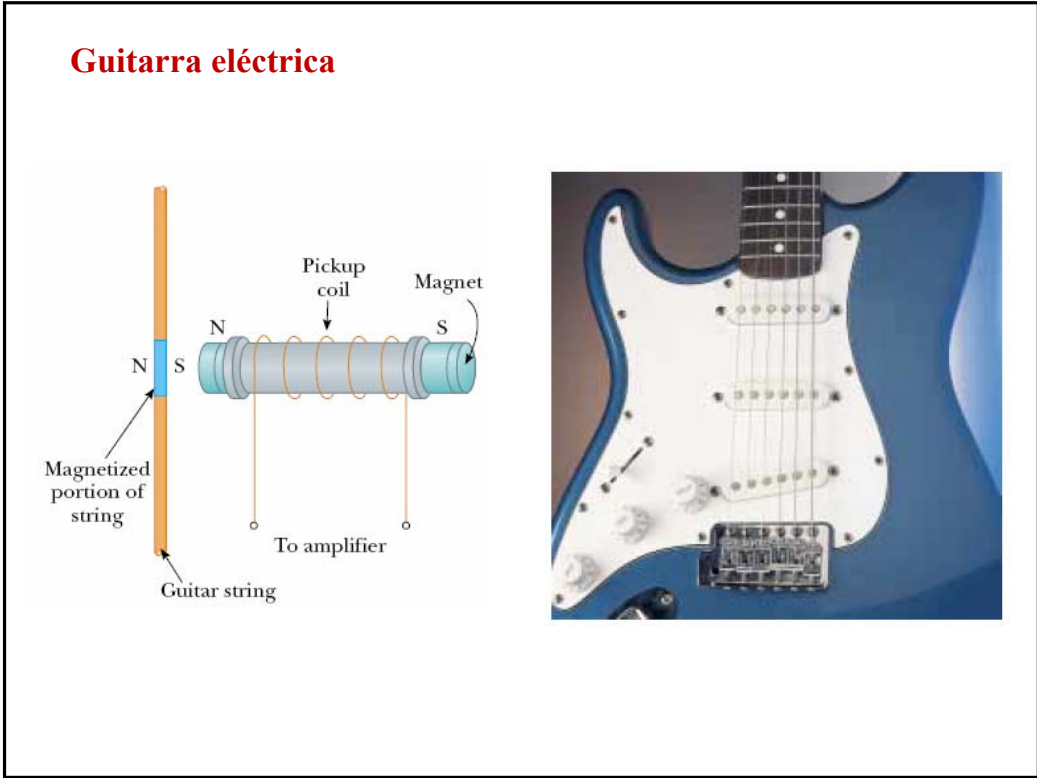
$$\mathcal{E} = -N \frac{d\Phi_B}{dt} = -NAB \frac{d}{dt} (\cos \omega t) = NAB\omega \sin \omega t$$

$$\mathcal{E}_{\max} = NAB\omega$$

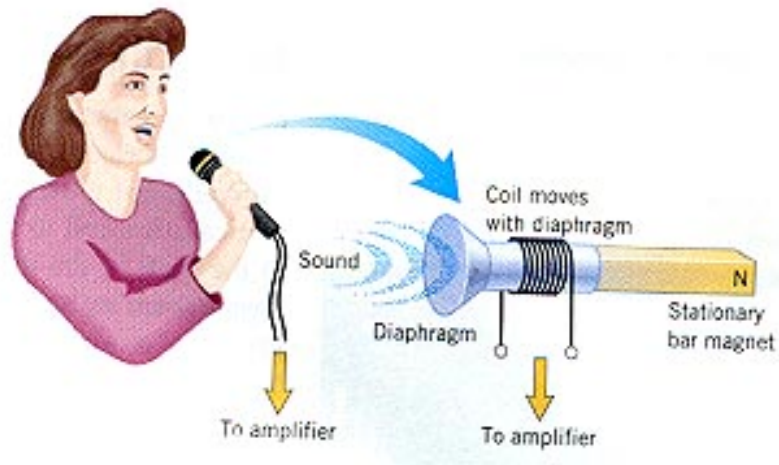




Guitarra eléctrica



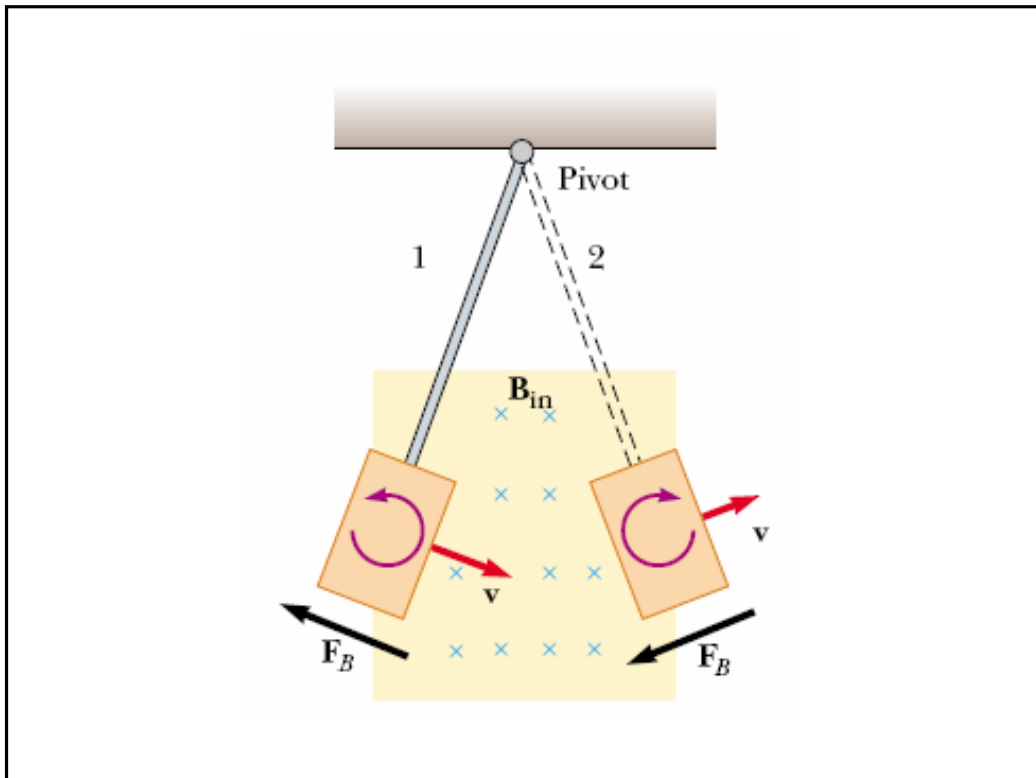
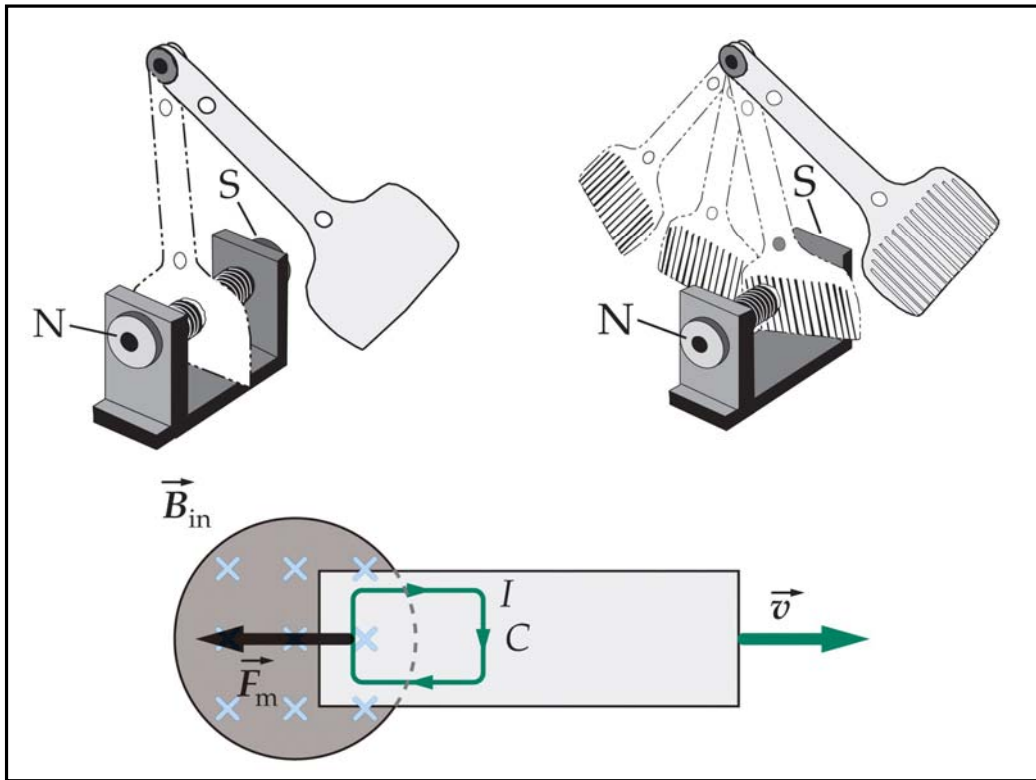
Micrófono de bobina móvil

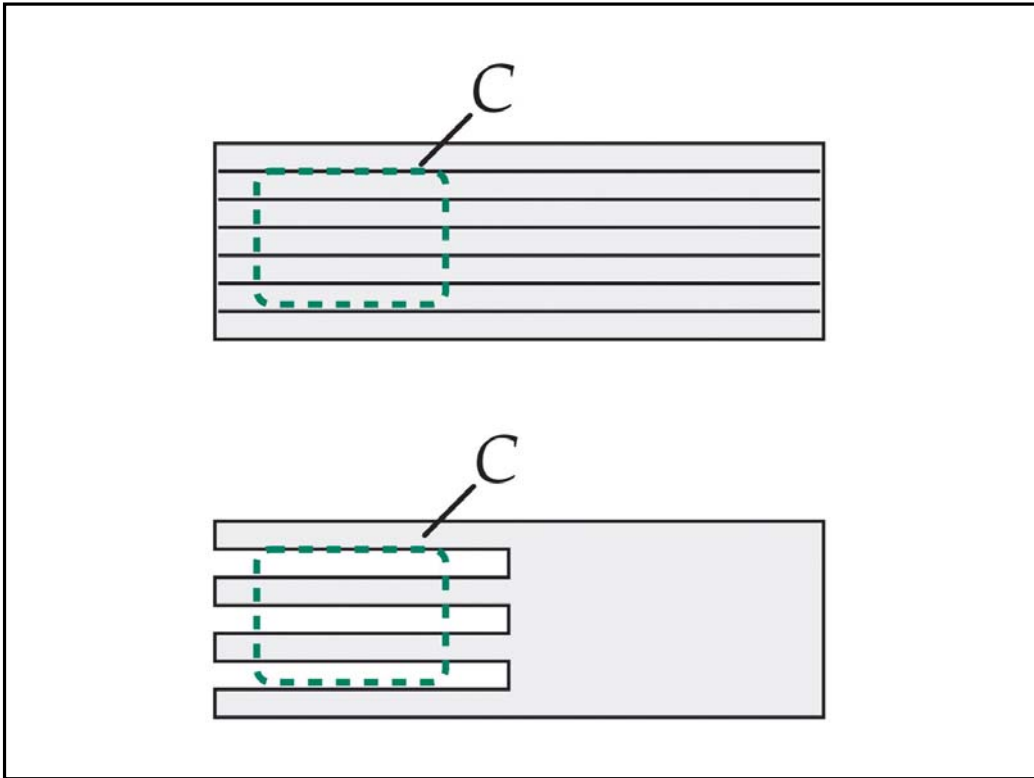


Corrientes de Foucault



Jean Leon Foucault (1819-1868)

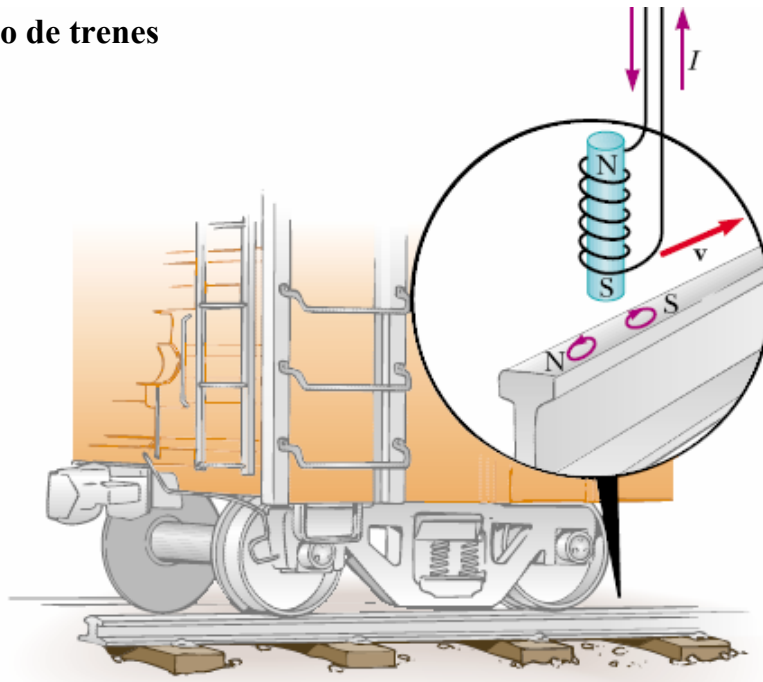




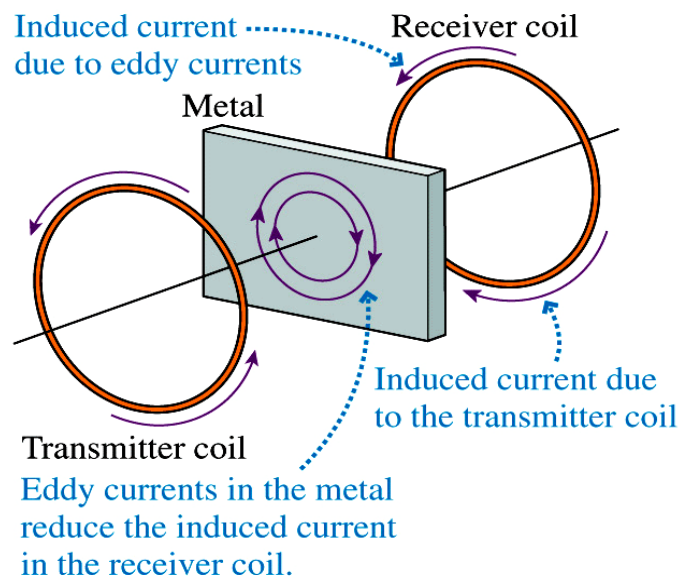
**Son útiles para
amortiguar
oscilaciones
mecánicas**

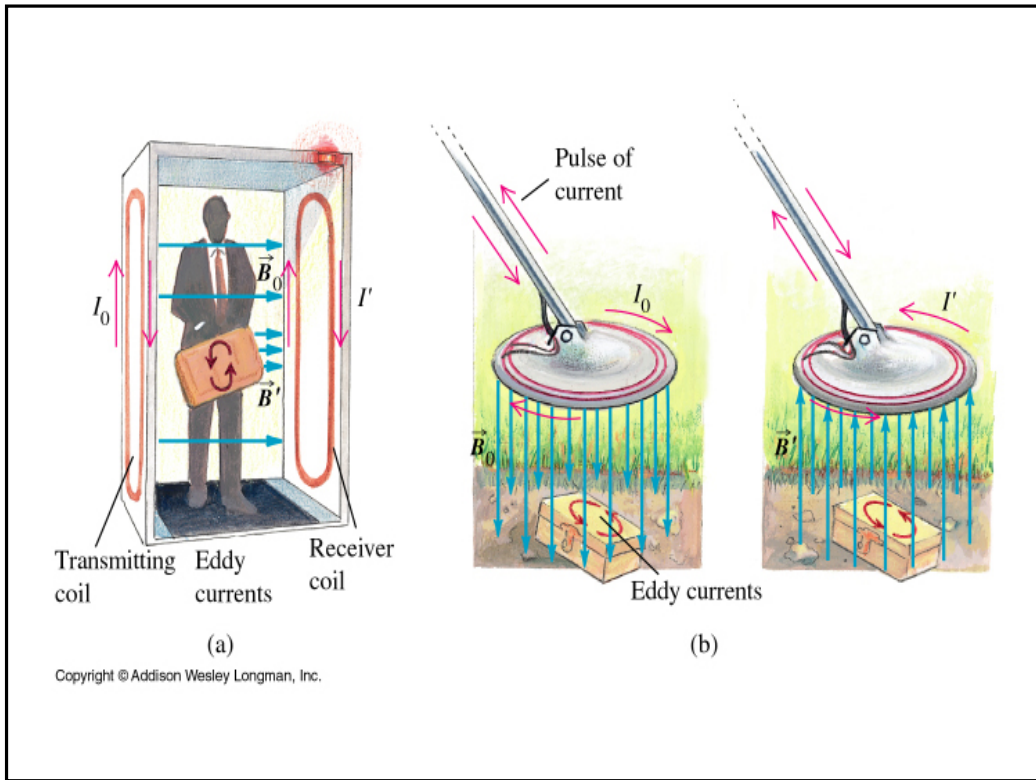


Freno de trenes

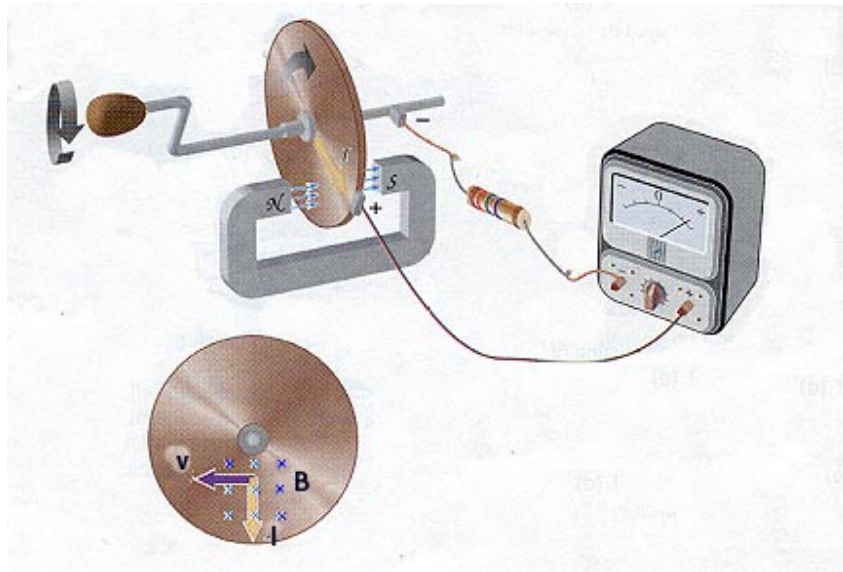


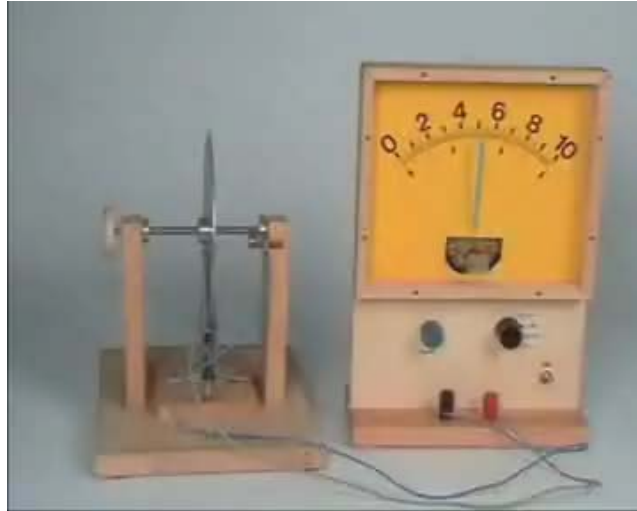
Detector de metales



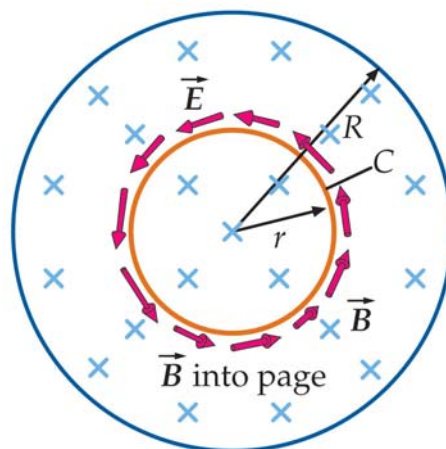


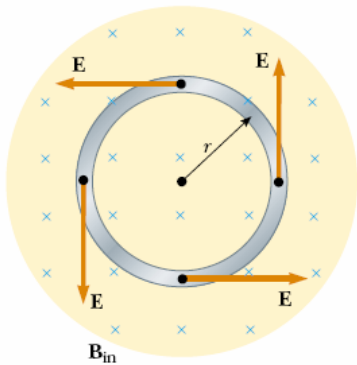
El primer generador





fem inducidas y campos eléctricos





$$\mathcal{E} = -\frac{d\phi_m}{dt}$$

$$\phi_m = \int_S \vec{B} \cdot \hat{n} dA = \int_S B_n dA$$

$$\mathcal{E} = \oint_C \vec{E}_{nc} \cdot d\vec{\ell} = -\frac{d}{dt} \int_S \vec{B} \cdot \hat{n} dA = -\frac{d\phi_m}{dt}$$

El campo eléctrico inducido no es conservativo

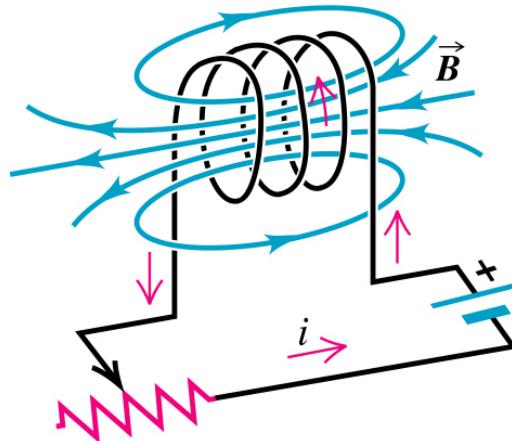
$$\oint_C \vec{E} \cdot d\vec{\ell} = -\frac{d}{dt} \int_S \vec{B} \cdot d\vec{A}$$

$$\left. \begin{aligned} \oint_L \vec{E} \cdot d\vec{\ell} &= \int_S \nabla \wedge \vec{E} \cdot d\vec{s} \\ \frac{d}{dt} \int_S \vec{B} \cdot d\vec{s} &= \int_S \frac{\partial \vec{B}}{\partial t} \cdot d\vec{s} \end{aligned} \right\} \rightarrow \nabla \wedge \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

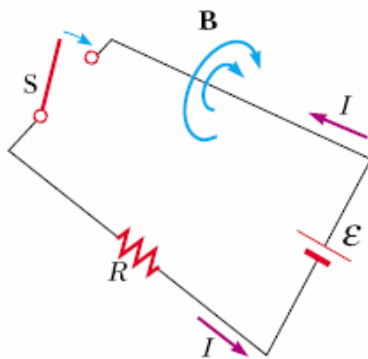
$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

La ley de Faraday-Henry en forma diferencial

Autoinducción



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S se cierra $\implies I(t)$

\Downarrow
B(t)

Variación temporal de Φ_B a través del área del circuito

fem autoinducida

$$\mathcal{E} = -\frac{d\Phi_B}{dt}$$

$$B \propto I$$

$$\Phi_B \propto I$$

Definimos:

$$\Phi_B = L I$$

L: coeficiente de autoinductancia o inductancia \implies

$$\mathcal{E}_L = -L \frac{dI}{dt}$$

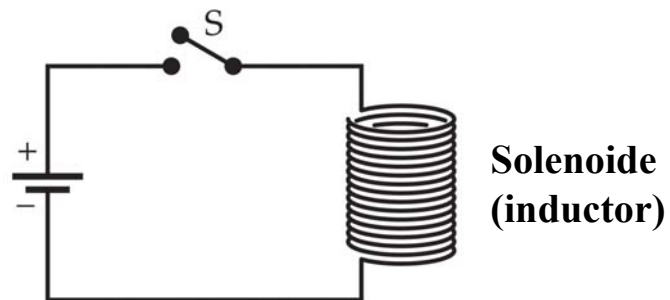
$$L = \Phi_B / I = -\mathcal{E}_L / (dI/dt)$$



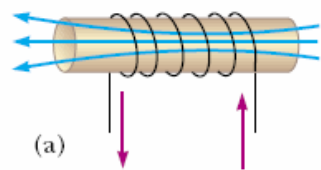
L

$$[L] = \text{T m}^2 / \text{A} = \text{Wb} / \text{A} = \text{V s} / \text{A} = \text{H (Henry)}$$

La inductancia de un circuito depende de su geometría

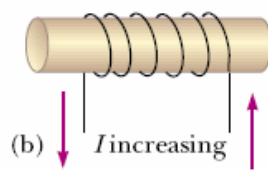


B



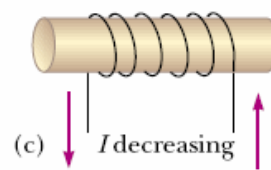
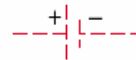
(a)

Lenz's law emf



(b) *I* increasing

Lenz's law emf



(c) *I* decreasing

$$L = \frac{N\Phi_B}{I}$$

$$B = \mu_0 n I = \mu_0 \frac{N}{\ell} I$$

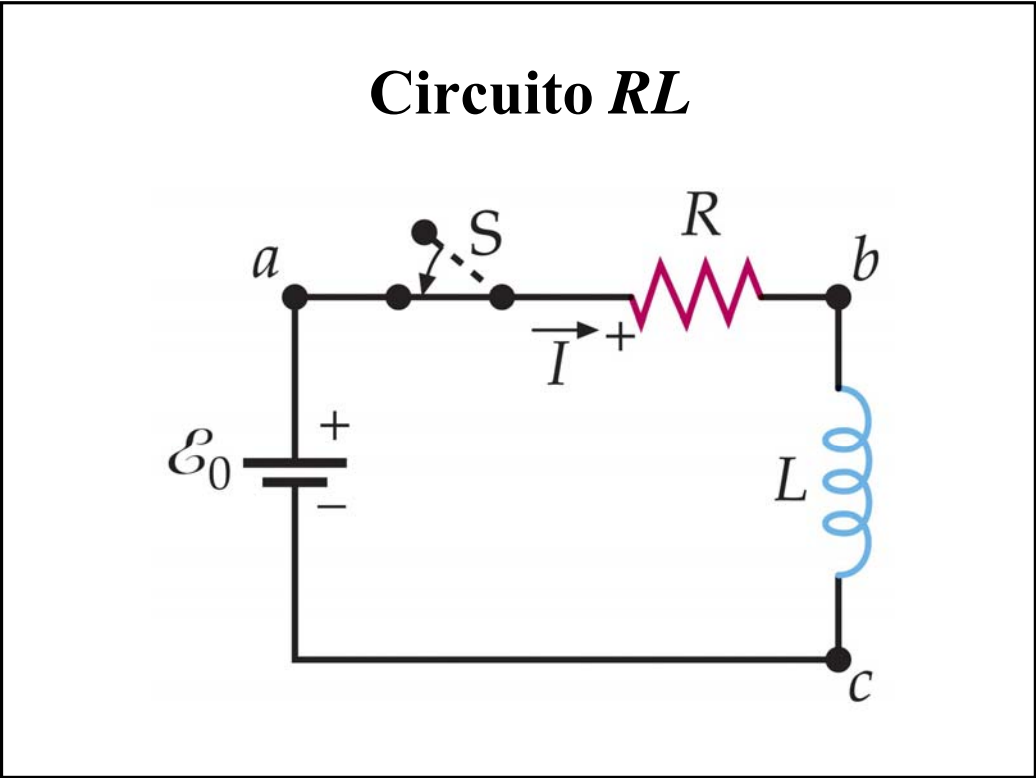
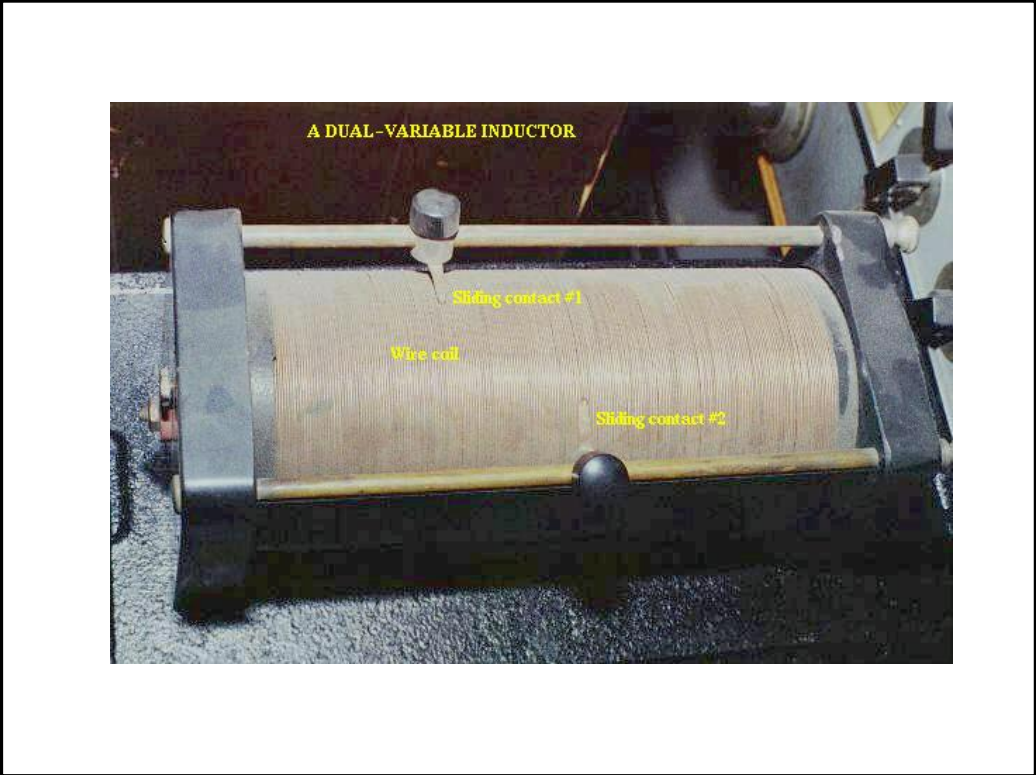
$$n = N/\ell$$

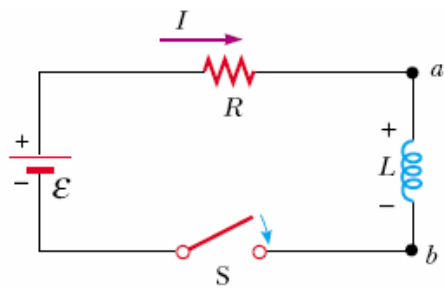
$$\Phi_B = BA = \mu_0 \frac{NA}{\ell} I$$

$$L = \frac{N\Phi_B}{I} = \frac{\mu_0 N^2 A}{\ell}$$

$$L = \mu_0 \frac{(n\ell)^2}{\ell} A = \mu_0 n^2 A \ell = \mu_0 n^2 V$$

$$V = A\ell$$





$$\mathcal{E} - IR - L \frac{dI}{dt} = 0$$

$$x = (\mathcal{E}/R) - I, \quad dx = -dI.$$

$$x + \frac{L}{R} \frac{dx}{dt} = 0 \quad \frac{dx}{x} = -\frac{R}{L} dt \quad \int_{x_0}^x \frac{dx}{x} = -\frac{R}{L} \int_0^t dt$$

$$\ln \frac{x}{x_0} = -\frac{R}{L} t \quad \Longrightarrow \quad x = x_0 e^{-Rt/L}$$

$$x_0 = x(t=0) = \mathcal{E}/R$$

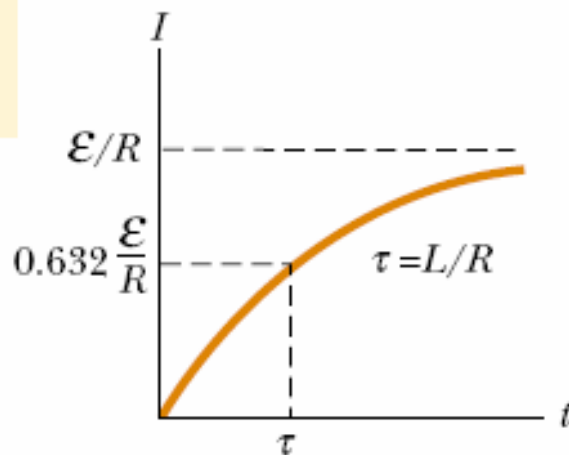
$$\frac{\mathcal{E}}{R} - I = \frac{\mathcal{E}}{R} e^{-Rt/L}$$

$$I = \frac{\mathcal{E}}{R} (1 - e^{-Rt/L})$$

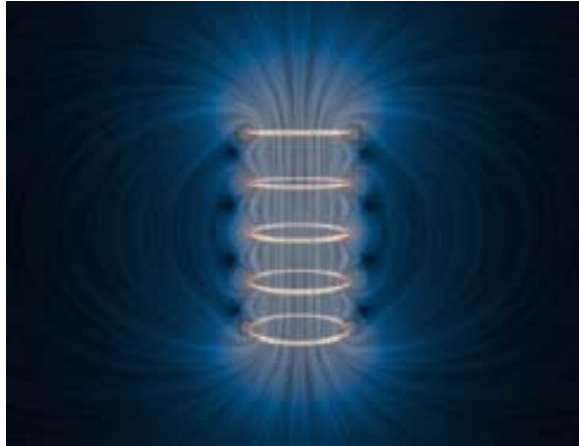
$$I = \frac{\mathcal{E}}{R} (1 - e^{-t/\tau})$$

$$\tau = \frac{L}{R}$$

**Constante de tiempo
del circuito RL**



Energía del Campo Magnético



La fem inducida evita que la batería establezca una corriente: la batería efectúa trabajo contra el inductor

Energía suministrada por la batería \longrightarrow Calor Joule + Energía almacenada en el inductor

$$\mathcal{E} - IR - L \frac{dI}{dt} = 0$$

$$I\mathcal{E} = I^2R + LI \frac{dI}{dt}$$

U: energía almacenada en el inductor

$$\frac{dU}{dt} = LI \frac{dI}{dt}$$

$$U = \int dU = \int_0^I LI dI = L \int_0^I I dI$$

$$U = \frac{1}{2} LI^2$$

$$U = \frac{1}{2} C(\Delta V)^2.$$

Análogo eléctrico

Para calcular la densidad de energía almacenada en el campo magnético, consideramos un solenoide:

$$L = \mu_0 n^2 A \ell$$

El campo magnético dentro del solenoide es $B = \mu_0 n I$

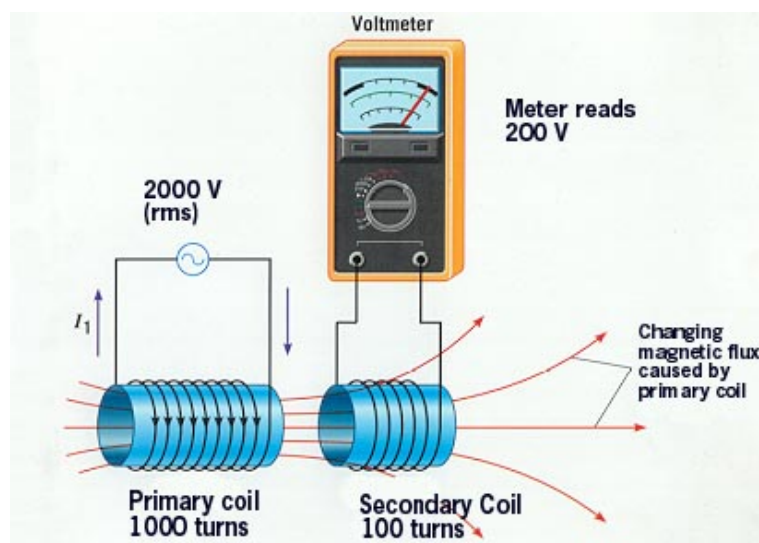
$$U = \frac{1}{2} L I^2 = \frac{1}{2} \mu_0 n^2 A \ell \left(\frac{B}{\mu_0 n} \right)^2 = \frac{B^2}{2\mu_0} A \ell$$

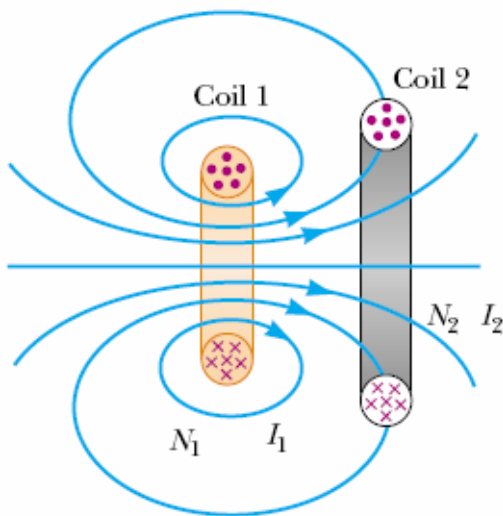
$A \ell$ volumen del solenoide donde se encuentra confinado B

$$u_B = \frac{U}{A \ell} = \frac{B^2}{2\mu_0}$$

Es válida para cualquier región del espacio donde haya B

Inductancia mutua





Φ_{12} .

Flujo magnético a través de la bobina 2 producido por la bobina 1.

Definimos inductancia mutua de la bobina 2 respecto a la 1

$$M_{12} \equiv \frac{N_2 \Phi_{12}}{I_1}$$

M12 depende de la geometría de ambos circuitos y de la orientación de uno respecto del otro.

La fem inducida en la bobina 2 es:

$$\mathcal{E}_2 = -N_2 \frac{d\Phi_{12}}{dt} = -N_2 \frac{d}{dt} \left(\frac{M_{12} I_1}{N_2} \right) = -M_{12} \frac{dI_1}{dt}$$

De la misma manera:

$$\mathcal{E}_1 = -M_{21} \frac{dI_2}{dt}$$

Se puede demostrar que $M_{12} = M_{21} = M$

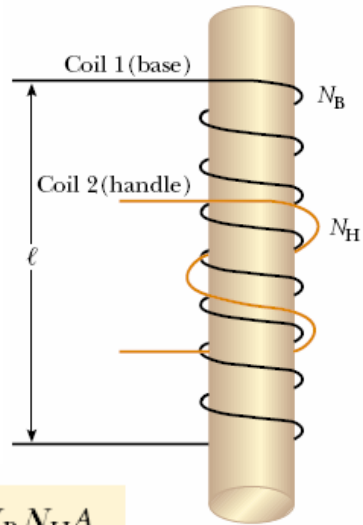
$$\mathcal{E}_2 = -M \frac{dI_1}{dt}$$

$$\mathcal{E}_1 = -M \frac{dI_2}{dt}$$

[M] = Hy

Ejemplo:

$$B = \frac{\mu_0 N_B I}{\ell}$$



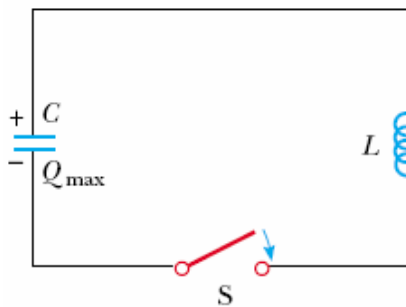
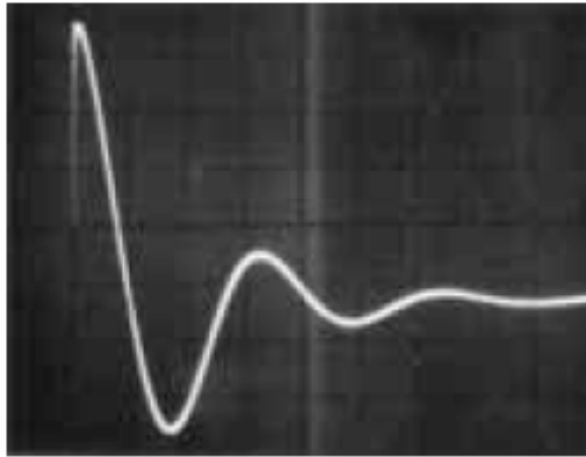
$$M = \frac{N_H \Phi_{BH}}{I} = \frac{N_H B A}{I} = \mu_0 \frac{N_B N_H A}{\ell}$$

Usualmente M se determina experimentalmente.



Oscilaciones eléctricas

Circuitos LC y RLC



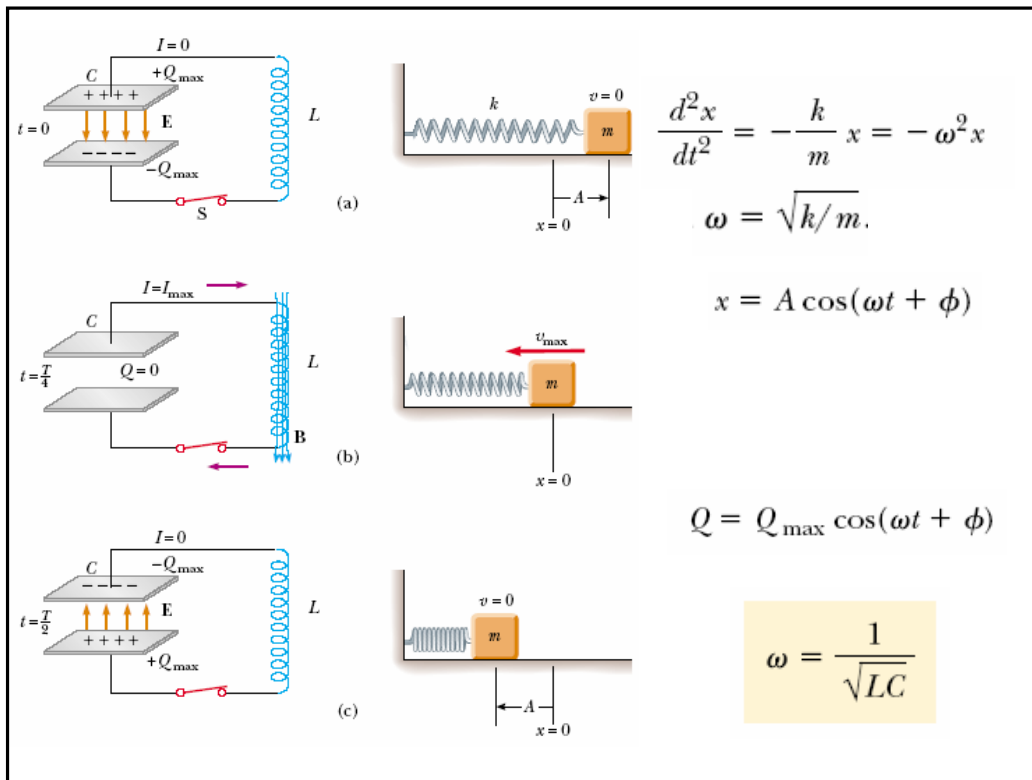
$$U = U_C + U_L = \frac{Q^2}{2C} + \frac{1}{2}LI^2$$

$$\frac{dU}{dt} = \frac{d}{dt} \left(\frac{Q^2}{2C} + \frac{1}{2}LI^2 \right) = \frac{Q}{C} \frac{dQ}{dt} + LI \frac{dI}{dt} = 0$$

$$I = dQ/dt. \longrightarrow \frac{Q}{C} + L \frac{d^2Q}{dt^2} = 0$$

$$\frac{d^2Q}{dt^2} = -\frac{1}{LC} Q$$

Ecuación diferencial del oscilador armónico

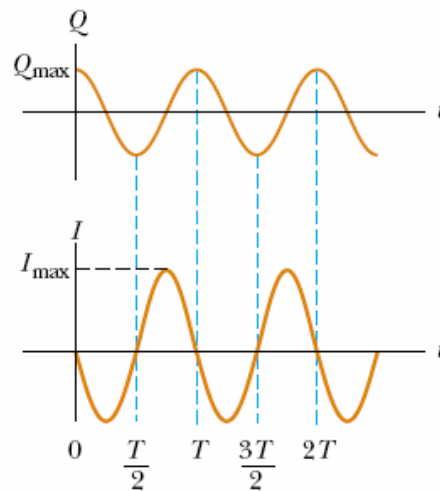


$$Q = Q_{\max} \cos(\omega t + \phi) \quad I = \frac{dQ}{dt} = -\omega Q_{\max} \sin(\omega t + \phi)$$

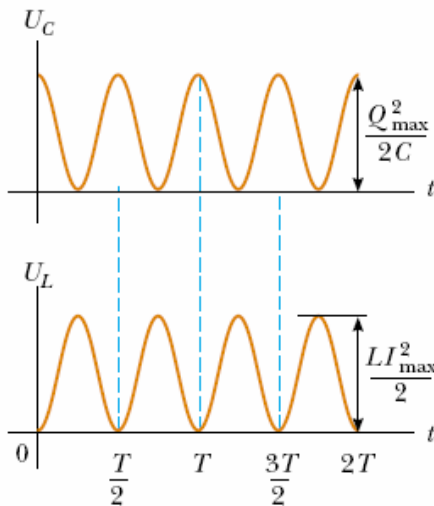
$$I = 0 \text{ at } t = 0 \quad 0 = -\omega Q_{\max} \sin \phi$$

$$Q = Q_{\max} \cos \omega t$$

$$I = -\omega Q_{\max} \sin \omega t = -I_{\max} \sin \omega t$$



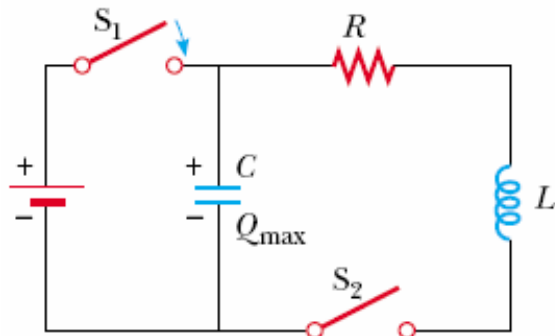
$$U = U_C + U_L = \frac{Q_{\max}^2}{2C} \cos^2 \omega t + \frac{1}{2} LI_{\max}^2 \sin^2 \omega t$$



$$\frac{Q_{\max}^2}{2C} = \frac{LI_{\max}^2}{2}$$

$$U = \frac{Q_{\max}^2}{2C} (\cos^2 \omega t + \sin^2 \omega t) = \frac{Q_{\max}^2}{2C}$$

**Que ocurre si
agrego R??**



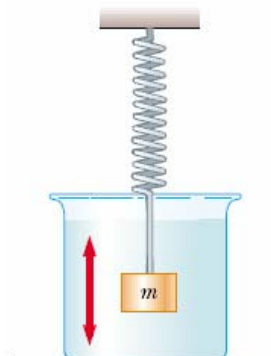
$$\frac{dU}{dt} = -I^2 R$$

$$LI \frac{dI}{dt} + \frac{Q}{C} \frac{dQ}{dt} = -I^2 R$$

$$LI \frac{d^2 Q}{dt^2} + I^2 R + \frac{Q}{C} I = 0$$

$$L \frac{d^2 Q}{dt^2} + IR + \frac{Q}{C} = 0$$

$$L \frac{d^2 Q}{dt^2} + R \frac{dQ}{dt} + \frac{Q}{C} = 0$$



$$L \frac{d^2 Q}{dt^2} + R \frac{dQ}{dt} + \frac{Q}{C} = 0 \leftrightarrow m \frac{d^2 x}{dt^2} + b \frac{dx}{dt} + kx = 0$$

$$\begin{aligned} Q &\leftrightarrow x \\ I &\leftrightarrow v_x \\ \Delta V &\leftrightarrow F_x \\ R &\leftrightarrow b \end{aligned}$$

$$m \frac{d^2 x}{dt^2} + b \frac{dx}{dt} + kx = 0$$

$$\begin{aligned} C &\leftrightarrow 1/k \\ L &\leftrightarrow m \end{aligned}$$

$$Q = Q_{\max} e^{-Rt/2L} \cos \omega_d t$$

donde

$$\omega_d = \left[\frac{1}{LC} - \left(\frac{R}{2L} \right)^2 \right]^{1/2}$$

$$I = \frac{dQ}{dt} \leftrightarrow v_x = \frac{dx}{dt}$$

$$\frac{dI}{dt} = \frac{d^2 Q}{dt^2} \leftrightarrow a_x = \frac{dv_x}{dt} = \frac{d^2 x}{dt^2}$$

$$U_L = \frac{1}{2} LI^2 \leftrightarrow K = \frac{1}{2} mv^2$$

$$U_C = \frac{1}{2} \frac{Q^2}{C} \leftrightarrow U = \frac{1}{2} kx^2$$

$$I^2 R \leftrightarrow bv^2$$

